

Rigid Body Mechanics

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Torque & Couples (HL)

Torque & Couples

Moment of a Force

- A moment is the **turning effect** of a force around a particular point
- Moments occur when forces cause objects to **rotate** about some pivot
- The moment of a force is given by

Moment (N m) = Force (N) × perpendicular distance from the pivot (m)

The SI unit for the moment is newton metres (N m)
 Moments of perpendicular & non-perpendicular forces



the perpendicular direction

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Your notes

Couples

- A couple is a **pair** of equal and opposite coplanar forces that act to produce **rotation** only
- A couple consists of a pair of forces that are:
 - Equal in magnitude
 - Opposite in direction
 - Perpendicular to the distance between them

Force couple diagram



- Unlike moments of a single force, the moment of a couple doesn't depend on a pivot
- The moment of a couple is equal to:

Moment (N m) = Force $(N) \times$ Perpendicular distance between the lines of action of the forces (m)

- A couple does not produce a net linear force
 - However, it does produce a turning effect called a torque

Torque

- The change in rotational motion due to a turning force is called **torque**
- The torque of a force *F* about an axis is given by

$$\tau = Fr\sin\theta$$

Torque of a non-perpendicular force

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- τ = torque (N m)
- F = applied force (N)
- r = perpendicular distance between the axis of rotation and the line of action of the force (m)
- θ = angle between the force and the axis of rotation (°)

Torque of a couple on a steering wheel

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forces



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- **B** is incorrect as the forces are in the same direction
- C is incorrect as the forces are different in size
- D is incorrect as the distance between the forces is not perpendicular

Worked example

A ruler of length 0.3 m is pivoted at its centre.

Two equal and opposite forces of magnitude 4 N are applied to the ends of the ruler, creating a couple as shown in the diagram.



Determine the magnitude of the torque of the couple on the ruler.

Answer:



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Worked example

The forces acting on a bicycle pedal at different positions during a ride are shown in the diagram below.

The distance from the pedal to the axis of rotation is 24 cm.





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- In position **A**: When $\theta = 180^\circ$, sin $\theta = 0$, so $\tau = 0$ i.e. the force exerts no torque on the pedal. When the pedal is at the bottom, no amount of pushing down will produce any torque on the pedal.
- In position B: When the angle θ is 90° and the force F is 90 N, the torque τ = 0.24 × 90 × sin 90° = 21.6 N m
- In position C: When the angle θ is 60° and the force F is 75 N, the torque $\tau = 0.24 \times 75 \times \sin 60^\circ =$ **15.6 N m**
- In position **D**: When the angle θ is 30° and the force *F* is 150 N, the torque $\tau = 0.24 \times 150 \times \sin 30^\circ =$ **18.0 N m**



😧 Examiner Tip

The terminology in this section can get confusing. For example, a moment is not a 'turning force' - the turning force is only part of the moment, the moment is the **effect** that the turning force has on the system when applied at a **distance** from a turning point, or pivot.

Also, torque is another term for 'moment', but in mechanics, torque can be used to describe a couple as a specific case of a moment.

Ultimately, when you carry out calculations, make sure you can identify

- The magnitude of the **applied force**
- The **perpendicular distance** between the force and the turning point (along the line of action) Make sure you understand that the forces that make up a **couple** cannot share the same line of action (the line through the point at which the force is applied). An example of this is shown in the diagram below:





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Rotational Equilibrium (HL)

Rotational Equilibrium

- A system is said to be in rotational equilibrium when
 - There is **no** resultant force acting
 - There is **no** resultant torque acting
- An object in rotational equilibrium will therefore remain at rest, or rotate with a constant angular velocity
 - This is analogous to Newton's First Law for translational equilibrium
- This means a body is in rotational equilibrium if

The sum of the clockwise moments is equal to the sum of the anticlockwise moments

- This is also known as the principle of moments and can be applied to a range of scenarios, such as a balanced beam
 - A beam is an example of a rigid, extended body

A balanced beam in rotational equilibrium



When the resultant force and resultant torque are both zero, the beam will be in rotational equilibrium



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Your notes

Worked example

Four beams of the same length each have three forces acting on them.

Which of the beams is in rotational equilibrium?



In rotational equilibrium:

Total clockwise torque = Total anticlockwise torque

• Consider beam **C**, taking torques from the centre of the beam (where its weight acts) :

Torque,
$$\tau = Fr(as \sin 90^\circ = 1)$$

Total clockwise torque = 15 × 50 = 750 N cm

Total anticlockwise torque = 25 × 30 = 750 N cm

• The total clockwise torque (750 N cm) = total anticlockwise torque (750 N cm), therefore, beam **C** is in rotational equilibrium

The other beams are **not** in rotational equilibrium because...

- Beam **A** has a resultant torque of 310 N cm anticlockwise
- Beam **B** has a resultant torque of 370 N cm clockwise
- Beam **D** has a resultant torque of 1790 N cm clockwise

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Examiner Tip

When considering an object in rotational equilibrium, choosing certain points can simplify calculations of resultant torque. Remember you can choose **any** point, not just the axis of rotation.

To simplify your calculation, choose a point where the torque of (most of) the forces are unknown, or when you need to determine where the resultant torque is **zero**. To do this, choose a point through which the lines of action of the forces pass



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Your notes

Unbalanced Torque

- In the same way that a resultant force produces linear acceleration, a **resultant torque** produces **angular acceleration**
- The direction of the angular acceleration depends on the direction of the net resultant torque Beam with an unbalanced torque



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If there is a net resultant torque in the clockwise or anti-clockwise direction, the beam will also have an angular acceleration in that direction

Worked example

A uniform plank of mass 30 kg and length 10 m is supported at its left end and at a point 1.5 m from the centre.





Calculate the maximum distance r to which a boy of mass 50 kg can walk without tipping the rod over.

Answer:

Step 1: Analyse the scenario and identify the forces

- Let the forces at each support be F_L (reaction force from the left support) and F_R (reaction force from the right support)
 - These are vertically upwards
- Just before the plank tips over, the system is in rotational equilibrium
- When the plank begins to tip over, the left support force F_L will become zero since the rod will no longer touch the support





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Your notes



Direction of angular acceleration: **clockwise**



You should know that torque is a vector quantity, however, at this level, you will only need to consider whether it produces clockwise or anti-clockwise motion



Your notes

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Angular Displacement, Velocity & Acceleration (HL)

Angular Displacement, Velocity & Acceleration

Angular displacement

• A rigid rotating body can be described using the following properties:

- Where:



Angular Velocity

- The angular velocity ω of a rigid rotating body is defined as:
 The rate of change in angular displacement with respect to time
- Angular velocity is measured in rad s⁻¹
- This can be expressed as an equation:

$$\omega = \frac{\Delta \theta}{\Delta t}$$

- Where:
 - ω = angular velocity (rad s⁻¹)
 - $\Delta \theta$ = angular displacement (rad)

• $\Delta t = change in time (s)$

Angular velocity to linear velocity

• The linear speed v is related to the angular speed ω by the equation:

 $v = r\omega$

- Where:
 - v = linear speed (m s⁻¹)
 - *r* = distance from the axis of rotation (m)
- Taking the angular displacement of a complete cycle as 2π, angular velocity ω can also be expressed as:

$$\omega = \frac{v}{r} = 2\pi f = \frac{2\pi}{T}$$

Rearranging gives the expression for linear speed:

$$v = 2\pi fr = \frac{2\pi r}{T}$$

- Where:
 - *f* = frequency of the rotation (Hz)
 - T = time period of the rotation (s)

Angular Acceleration

• Angular acceleration α is defined as

The rate of change of angular velocity with time

- Angular acceleration is measured in rad s⁻²
- This can be expressed as an equation:



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- Where:
 - α = angular acceleration (rad s⁻²)
 - $\Delta \omega$ = change in angular velocity, or $\Delta \omega = \omega_f \omega_i$ (rad s⁻¹)
 - Δt = change in time (s)

Angular acceleration to linear acceleration

• Using the definition of angular velocity ω with the equation for angular acceleration α gives:

$$\Delta \omega = \frac{\Delta v}{r}$$
$$\alpha = \frac{\Delta \omega}{\Delta t} = \frac{\Delta v}{r\Delta t} = \frac{a}{r}$$

• Rearranging gives the expression for linear acceleration:

$$a = r\alpha$$

- Where:
 - $a = \text{linear acceleration (m s}^{-2})$
 - r = distance from the axis of rotation (m)
 - ΔV = change in linear velocity, or $\Delta V = V u$ (m s⁻¹)

Graphs of Rotational Motion

• Graphs of rotational motion can be interpreted in the same way as linear motion graphs



Your notes



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displacement	$s = r\theta$	$\theta = \frac{s}{r}$
velocity	$V = T \omega$	$\omega = \frac{v}{r}$
acceleration	$a = r\alpha$	



Examiner Tip

While there are many similarities between the angular quantities used in this topic and the angular quantities used in the circular motion topic, make sure you are clear on the distinctions between the two, for example, **angular** acceleration and **centripetal** acceleration are **not** the same thing!

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Angular Acceleration Formula (HL)

Angular Acceleration Formula

- The kinematic equations of motion for uniform linear acceleration can also be re-written for rotational motion
- The four kinematic equations for uniform **linear** acceleration are

$$v = u + at$$

$$s = ut + \frac{1}{2}at^{2}$$

$$v^{2} = u^{2} + 2as$$

$$s = \frac{(u + v)t}{2}$$

• This leads to the four kinematic equations for uniform **rotational** acceleration

$$\omega_{i} = \omega_{f} + \alpha t$$
$$\Delta \theta = \omega_{i} t + \frac{1}{2} \alpha t^{2}$$
$$\omega_{f}^{2} = \omega_{i}^{2} + 2\alpha \Delta \theta$$
$$\Delta \theta = \frac{(\omega_{i} + \omega_{f})t}{2}$$

• The five linear variables have been swapped for the rotational equivalents, as shown in the table below

Variable	Linear	Rotational
displacement	S	θ
initial velocity	u	$\omega_{ m i}$
final velocity	V	ω_{f}
acceleration	а	α







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time	+	+
time	l	L



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Worked example

The turntable of a record player is spinning at an angular velocity of 45 RPM just before it is turned off. It then decelerates at a constant rate of 0.8 rad s⁻².

Determine the number of rotations the turntable completes before coming to a stop.

Answer:

Step 1: List the known quantities

- Initial angular velocity, $\omega_i = 45 \text{ RPM}$
- Final angular velocity, $\omega_f = 0$
- Angular acceleration, α = 0.8 rad s⁻²
- Angular displacement, $\Delta \theta = ?$

Step 2: Convert the angular velocity from RPM to rad $\ensuremath{\mathsf{s}}^{-1}$

• One revolution corresponds to 2π radians, and RPM = revolutions per minute, so

$$\omega = 2\pi f$$
 and $f = \frac{RPM}{60}$ (to convert to seconds)

$$\omega_i = \frac{2\pi \times \text{RPM}}{60} = \frac{2\pi \times 45}{60} = \frac{3\pi}{2} \text{ rad s}^{-1}$$

Step 3: Select the most appropriate kinematic equation

• We know the values of ω_i , ω_f and α , and we are looking for angular displacement θ , so the best equation to use would be

$$\omega_f^2 = \omega_i^2 + 2\alpha \Delta \theta$$

Step 4: Rearrange and calculate the angular displacement $\Delta heta$

$$0 = \omega_i^2 - 2\alpha \Delta \theta$$

$$\Delta \theta = \frac{\omega_i^2}{2\alpha} = \frac{\left(\frac{3\pi}{2}\right)^2}{2 \times 0.8}$$

Angular displacement, $\Delta \theta$ = 13.88 rad

Step 5: Determine the number of rotations in $\Delta \, heta$

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- There are 2π radians in 1 rotation
- Therefore, the number of rotations = $\frac{13.88}{2\pi}$ = 2.2

Your notes

• This means the turntable spins 2.2 times before coming to a stop

Moment of Inertia (HL)

Moment of Inertia

- In linear motion, the resistance to a change of motion, i.e. linear acceleration, is known as inertia
 The larger the mass an object has, the greater its inertia
- In rotational motion, the distribution of mass around an axis must be considered, using moments of inertia
- The moment of inertia of a rigid, extended body is defined as:

The resistance to a change of rotational motion, depending on the distribution of mass around a chosen axis of rotation

- Moment of inertia is measured in kg m²
- The moment of inertia of a body corresponds to how 'easy' or 'hard' it is to rotate, and this is dependent on many factors, including
 - Its shape
 - Its density
 - Its orientation (relative to an axis of rotation)
- These factors allow an object's **distribution of mass** to be taken into account
 - It also means that the moment of inertia of a singular object can **change** depending on its orientation in relation to the chosen axis of rotation
- For example, the moment of inertia of a **thin rod** is different for each of the following orientations:
 - Rotation about its **vertical axis**
 - Rotation about its centre of mass
 - Rotation about one end





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The moment of inertia of a body can change depending on its orientation relative to the axis of rotation

- These are just a few of the possible orientations of the axis of rotation for a thin rod
 - There is an **infinite** range of possible axes, and therefore moments of inertia
 - This also applies to nearly all rigid, extended objects that could be considered

Examiner Tip

Make sure you are clear on the distinction between linear motion and rotational motion here. The implications of considering the distribution of masses in relation to an axis of rotation, as opposed to considering them as uniform, have important consequences when carrying out calculations



Calculating Moments of Inertia

• The moment of inertia I of a point mass is equal to

 $I = mr^2$

- Where:
 - $I = \text{moment of inertia} (\text{kg m}^2)$
 - m = mass of the object (kg)
 - *r* = distance from its axis of rotation (m)
- A point mass is the simplest type of object to consider, the moment of inertia of a non-point mass can be calculated using

$$I = \sum mr^2$$

- This means that the sum of the moments of inertia of all the point masses in the system gives the total moment of inertia of the system
- Some moments of inertia of common shapes are shown below:



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Worked example

Two solid spheres form a dumbbell when attached to each end of a thin rod. The dumbbell rotates with the centre of mass of each sphere at a distance of 22 cm from the axis of rotation, as shown in the diagram.

The thin rod has a mass of 20 g. Each sphere has a radius of 4 cm and a mass of 750 g.



Moment of inertia of a thin rod about its centre =
$$\frac{1}{12} mL^2$$

Moment of inertia of a solid sphere =
$$\frac{2}{5}mr^2$$

Determine

- (a) the overall moment of inertia of the dumbbell arrangement
- (b) the ratio of the moment of inertia of the thin rod to the overall moment of inertia of the dumbbell arrangement

Answer:

(a)

• The overall moment of inertia of the dumbbell is the sum of all the moments of inertia in the arrangement

$$I = \sum mr^2 = 2 \times \left(\frac{2}{5}m_{sphere}r^2\right) + \frac{1}{12}m_{rod}L^2$$

Where:

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- Mass of a sphere, $m_{sphere} = 750 \text{ g} = 0.75 \text{ kg}$
- Distance from axis to each sphere, *I* = 22 cm = 0.22 m
- Mass of the rod, *m*_{rod} = 20 g = 0.02 kg
- Length of the rod, $L = 2 \times (22 4) = 36 \text{ cm} = 0.36 \text{ m}$

$$I = 2 \times \left(\frac{2}{5} \times 0.75 \times 0.22^2\right) + \left(\frac{1}{12} \times 0.02 \times 0.36^2\right)$$

Moment of inertia of the dumbbell: $I = 0.029 \text{ kg m}^2$

(b)

• The moment of inertia of the thin rod is

$$I_{rod} = \frac{1}{12} \times 0.02 \times 0.36^2 = 2.16 \times 10^{-4} \,\mathrm{kg}\,\mathrm{m}^2$$

- Therefore, the ratio I_{rod} / I is

$$\frac{I_{rod}}{I} = \frac{2.16 \times 10^{-4}}{0.029} = 0.0071$$

• This means the rod contributes about 0.7% of the overall moment of inertia of the dumbbell

Examiner Tip

You will never be expected to memorise the moments of inertia of different shapes, they will **always** be given in an exam question where required



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Newton's Second Law for Rotation (HL)

Newton's Second Law for Rotation

In linear motion, the **force** required to give an object a certain **acceleration** depends on its **mass**

F = ma

- This is Newton's Second Law of linear motion, where:
 - F = force(N)
 - m = mass (kg)
 - $a = \text{linear acceleration (m s}^{-2})$
- In rotational motion, the torque required to give a rotating object a certain angular acceleration depends on its moment of inertia

 $\tau = I\alpha$

- This is Newton's Second Law of **rotational motion**, where:
 - τ = torque (N m)
 - $I = \text{moment of inertia} (\text{kg m}^2)$
 - α = angular acceleration (rad s⁻²)



Newton's second law for rotating bodies is equivalent to Newton's second law for linear motion



• This equation comes from the fact that **torque** is the rotational equivalent of **force**:

Force:
$$F = ma$$

Torque:
$$\tau = Fr$$

- Where:
 - Γ = distance from the axis of rotation (m)
- Combining these equations gives:

$$\tau = r(ma)$$

- The moment of inertia of a rotating body can be thought of as analogous to mass
 - The inertia of a mass describes its ability to resist changes to **linear** motion, which is referring to **linear acceleration**
 - Similarly, the moment of inertia of a mass describes its ability to resist changes to **rotational** motion, which is referring to **angular acceleration**

Angular acceleration:
$$\alpha = \frac{a}{r}$$

Moment of inertia (point mass): $I = mr^2$

• Using these equations with the equations for force and torque leads to:

$$\tau = r(mr\alpha)$$
$$\tau = (mr^2)\alpha$$
$$\tau = I\alpha$$

Comparison of linear and rotational variables in Newton's Second Law

Linear variable	Rotational variable
Force, F	Torque, $ au$
Mass, <i>M</i>	Moment of inertia, I
Acceleration, a	Angular acceleration, $lpha$
Newton's Second Law, $F lpha a$	Newton's Second Law, $ au lpha lpha$
F = ma	$\tau = I\alpha$

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Worked example

A block of mass *m* is attached to a string that is wrapped around a cylindrical pulley of mass *M* and radius *R*, as shown in the diagram.

The moment of inertia of the cylindrical pulley about its axis is $\frac{1}{2}MR^2$.



When the block is released, the pulley begins to turn as the block falls.

Write an expression for the acceleration of the block.

Answer:

Step 1: Identify the forces acting on the block



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$$mg - T = ma \text{ eq. (1)}$$

Step 3: Apply Newton's second law to the rotation of the pulley

$$\tau = I\alpha$$
$$TR = I\alpha$$

Step 4: Write the equation for the pulley in terms of acceleration a

• The angular acceleration lpha of the pulley is:

$$\alpha = \frac{a}{R}$$

• Substitute this into the previous equation:

$$TR = I \frac{a}{R}$$

• Substitute in the expression for the moment of inertia and simplify:

Moment of inertia of the cylinder:
$$I = \frac{1}{2}MR^2$$

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$$TR = \left(\frac{1}{2}MR^2\right)\frac{a}{R}$$
$$T = \left(\frac{1}{2}MR^2\right)\frac{a}{R^2} = \frac{1}{2}Ma$$
$$T = \frac{1}{2}Ma \text{ eq. (2)}$$

Step 5: Substitute eq. (2) into eq. (1) and rearrange for acceleration a

$$mg - \frac{1}{2}Ma = ma$$

$$mg = ma + \frac{1}{2}Ma = a\left(m + \frac{M}{2}\right)$$
Acceleration of the block: $a = \frac{mg}{m + \frac{M}{2}}$

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Angular Momentum (HL)

Angular Momentum

Angular momentum is the rotational equivalent of linear momentum, which is defined by mass x velocity, or

$$p = mv$$

• Therefore, angular momentum L is defined by

$$L = I\omega$$

- Where:
 - L = angular momentum (kg m² rad s⁻¹)
 - $I = \text{moment of inertia} (\text{kg m}^2)$
 - ω = angular velocity (rad s⁻¹)

Angular Momentum of a Point Mass

• The moment of inertia of a rotating point mass *m* which is a distance *r* from an axis of rotation is equal to

$$I = mr^2$$

• The angular velocity of the point mass is given by

$$\omega = \frac{V}{r}$$

• Therefore, the angular momentum of the point mass is equal to

$$L = I\omega = (mr^2) \times \frac{v}{r} = mvr$$



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Worked example

A horizontal rigid bar is pivoted at its centre so that it is free to rotate. A point particle of mass 3*M* is attached at one end of the bar and a container is attached at the other end, both are at a distance of *R* from the central pivot.

A point particle of mass M moves with velocity v at right angles to the rod as shown in the diagram.



The particle collides with the container and stays within it as the system starts to rotate about the vertical axis with angular velocity ω .

The mass of the rod and the container are negligible.

Write an expression for the angular momentum of the system about the vertical axis:

- (a) just before the collision, in terms of *M*, *v* and *R*
- (b) just after the collision, in terms of M, R and ω .

Answer:

(a) Just before the collision:

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• Angular momentum is equal to:

 $L = I\omega$

• The moment of inertia of a point particle is

 $I = mr^2$

• Linear velocity is related to angular velocity by

 $v = \omega r$

- The rod, container and 3*M* mass are all stationary before the collision, so we only need to consider the angular momentum of the point particle
- Where:
 - Mass of the particle, m = M
 - Distance of the particle from the axis, r = R
 - Angular velocity of the particle, $\omega = \frac{V}{R}$
- Therefore, the angular momentum of the system before the collision is:

$$L = (MR^2) \times \frac{V}{R} = MVR$$

(b) After the collision:



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- The whole system rotates with an angular velocity of ω
- We are considering the rod and the container as massless, so we only need to consider the angular momentum of the two masses *M* and 3*M*
- Therefore, the angular momentum of the system after the collision is:

$$L = (MR^2)\omega + (3MR^2)\omega$$

$$L = 4MR^2\omega$$

Examiner Tip

You should know that objects travelling in straight lines **can** have angular momentum – just make sure you understand that it all depends on the position of the object in relation to the axis of rotation being considered





Conservation of Angular Momentum

- As with linear momentum, angular momentum is **always conserved**
- The principle of conservation of angular momentum states:

The angular momentum of a system always remains constant, unless a net torque is acting on the system

- This conservation law has many real-world applications, for example
 - A person on a **spinning chair** spins faster while their arms and legs are contracted and slower while extended
 - Objects in **elliptical orbits** travel faster nearer the object they orbit and slower when further away
 - Ice skaters can change their rotation velocity by extending or contracting their arms
 - Tornados spin faster as their radius decreases



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Ice skaters can change their moment of inertia by extending or contracting their arms and legs. Due to conservation of angular momentum, this allows them to spin faster or slower



Problems involving a **constant angular momentum** can be solved using the equation:

$$I_i \omega_i = I_f \omega_f$$

- Where:
 - I_i = initial moment of inertia (kg m²)
 - ω_i = initial angular velocity (rad s⁻¹)
 - I_f = final moment of inertia (kg m²)
 - \mathcal{O}_f = final angular velocity (rad s⁻¹)

Worked example

The diagram shows the different positions of a diver between jumping off a springboard and entering the water.

During their fall, the diver pulls their arms and legs into a tight tuck position while in the air and straightens them before entering the water.







Α.	increases	increases
В.	decreases	increases
C.	increases	decreases
D.	decreases	decreases

Answer: B

- After the diver leaves the springboard, there is no longer a resultant torque acting on them
 - This means their angular momentum remains **constant** throughout the dive
- Due to the conservation of angular momentum:

$$I_i \omega_i = I_f \omega_f = constant$$

- When the diver tucks their arms and legs in closer to their body, they **decrease** their moment of inertia
 - This eliminates options **A** & **C**
- Therefore, to conserve angular momentum, when the diver's moment of inertia decreases, their angular velocity must **increase**



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Worked example

A spherical star of mass M and radius R rotates about its axis. The star explodes, ejecting mass in space

radially and symmetrically. The remaining star is left with a mass of $\frac{1}{10}M$ and a radius of $\frac{1}{50}R$.

Calculate the ratio of the star's final angular velocity to its initial angular velocity.

The moment of inertia of a sphere is $rac{2}{5}MR^2$

Answer:

- Before the star explodes:
 - Initial moment of inertia, $I_i = \frac{2}{5}MR^2$
 - Initial angular velocity = ω_i
- After the star explodes:

Final moment of inertia,
$$I_f = \frac{2}{5} \left(\frac{1}{10} M\right) \left(\frac{1}{50} R\right)^2$$

- Final angular velocity = \mathcal{O}_{f}
- From the conservation of angular momentum:

$$\begin{pmatrix} \frac{2}{5}MR^2 \end{pmatrix} \omega_i = \left(\frac{2}{5} \times \frac{1}{10}M \times \left(\frac{1}{50}\right)^2 R^2\right) \omega_f$$
$$\begin{pmatrix} \frac{2}{5}MR^2 \end{pmatrix} \omega_i = \left(\frac{2}{5}MR^2\right) \frac{1}{25\,000}\,\omega_f$$

 $I.\omega = I.\omega_c$

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Therefore, the ratio of the star's final angular velocity to its initial angular velocity is:

$$\frac{\omega_f}{\omega_i} = 25\,000$$

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Angular Impulse (HL)

Angular Impulse

In linear motion, the resultant force on a body can be defined as the rate of change of linear momentum:

$$F = \frac{\Delta p}{\Delta t}$$

• This leads to the definition of linear impulse:

An average resultant force F acting for a time Δt produces a change in linear momentum Δp

$$\Delta p = F \Delta t = \Delta(mv)$$

• Similarly, the **resultant torque** on a body can be defined as the rate of change of angular momentum:

$$\tau = \frac{\Delta L}{\Delta t}$$

- Where:
 - τ = resultant torque on a body (N m)
 - ΔL = change in angular momentum (kg m² s⁻¹)
 - Δt = time interval (s)
- This leads to the definition of angular impulse:

An average resultant torque τ acting for a time $\Delta\,t$ produces a change in angular momentum $\Delta\,L$

$$\Delta L = \tau \Delta t = \Delta (I\omega)$$

- Angular impulse is measured in kg m² s⁻¹, or N m s
- This equation requires the use of a constant resultant torque
 - If the resultant torque changes, then an **average** of the values must be used
- Angular impulse describes the effect of a torque acting over a time interval
 - This means a small torque acting over a long time has the same effect as a large torque acting over a short time

Angular Impulse on a Torque-Time Graph

• The area under a torque-time graph is equal to the angular impulse or the change in angular momentum







When the torque is not constant, the angular impulse is the area under a torque-time graph

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Worked example

The graph shows the variation of time t with the net torque τ on an object which has a moment of inertia of 5.0 kg m².



At t = 0, the object rotates with an angular velocity of 2.0 rad s⁻¹ clockwise.

Determine the magnitude and direction of rotation of the angular velocity at t = 5 s.

In this question, take anticlockwise as the positive direction.

Answer:

Step 1: Use the graph to determine the angular impulse

• The area under a torque-time graph is equal to angular impulse, or the change in angular momentum

$$\Delta L = \tau \times \Delta t$$



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 ω_f = -1.0 rad s⁻¹ in the clockwise direction

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Rotational Kinetic Energy (HL)

Rotational Kinetic Energy

• A body moving with **linear velocity** has an associated linear kinetic energy given by

$$E_k = \frac{1}{2}mv^2$$

$$E_k = \frac{p^2}{2m}$$

Similarly, a rotating body with angular velocity has an associated rotational kinetic energy given by

$$E_k = \frac{1}{2}I\omega^2$$

$$E_k = \frac{L^2}{2I}$$

- Where:
 - E_{μ} = rotational kinetic energy (J)
 - $I = \text{moment of inertia} (\text{kg m}^2)$
 - ω = angular velocity (rad s⁻¹)
 - L = angular momentum (kg m² s⁻¹)

Rolling without slipping

- Circular objects, such as wheels, are made to move with **both** linear and rotational motion
 - For example, the wheels of a car, or bicycle **rotate** causing it to move **forward**
- Rolling motion without slipping is a combination of rotating and sliding (translational) motion
- When a disc rotates:
 - Each point on the disc has a **different** linear velocity depending on its distance from the centre $(V \propto I)$
 - The linear velocity is the **same** at all points on the **circumference**
- When a disc **slips**, or **slides**:
 - There is not enough friction present to allow the object to roll
 - Each point on the object has the **same** linear velocity
 - The angular velocity is **zero**
- So, when a disc **rolls without slipping**:
 - There is enough friction present to initiate rotational motion allowing the object to roll
 - The point in contact with the surface has a velocity of zero
 - The centre of mass has a velocity of $V = \omega r$

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Rolling motion is a combination of rotational and translational motion. The resultant velocity at the bottom is zero and the resultant velocity at the top is 2v

Rolling down a slope

- Another common scenario involving rotational and translational motion is an object (usually a ball or a disc) rolling down a slope
- At the top of the slope, a stationary object will have gravitational potential energy equal to

$$E_p = mg \Delta h$$

- As the object **rolls** down the slope, the gravitational potential energy will be transferred to **both** translational (linear) and rotational kinetic energy
- At the **bottom** of the slope, the **total kinetic energy** of the object will be equal to

$$E_{K \text{ total}} = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2$$

Your notes



- The linear or angular velocity can then be determined by
 - Equating E_p and E_K total
 - Using the equation for the moment of inertia of the object
 - Using the relationship between linear and angular velocity $V = \omega r$
- For example, for a ball (a solid sphere) of mass *m* and radius *r*, its moment of inertia is

$$I = \frac{2}{5}mr^2$$

- Equating the equations for ${\cal E}_p$ and ${\cal E}_K$ total and simplifying gives

$$mg\Delta h = \frac{1}{2}m(\omega r)^2 + \frac{1}{2}\left(\frac{2}{5}mr^2\right)\omega^2$$
$$mg\Delta h = \frac{1}{2}m\omega^2 r^2 + \frac{1}{5}m\omega^2 r^2$$

$$mg\Delta h = \frac{7}{10}m\omega^2 r^2$$

Worked example

A flywheel of mass M and radius R rotates at a constant angular velocity ω about an axis through its centre. The rotational kinetic energy of the flywheel is E_{K} .

The moment of inertia of the flywheel is $\frac{1}{2}MR^2$.

A second flywheel of mass $\frac{1}{2}M$ and radius $\frac{1}{2}R$ is placed on top of the first flywheel. The new angular velocity of the combined flywheels is $\frac{2}{3}\omega$.



What is the new rotational kinetic energy of the combined flywheels?

A.
$$\frac{E_K}{2}$$
 B. $\frac{E_K}{4}$ C. $\frac{E_K}{8}$ D. $\frac{E_K}{24}$

Answer: A

• The kinetic energy of the first flywheel is

$$E_{K} = \frac{1}{2}I\omega^{2} = \frac{1}{2} \times \left(\frac{1}{2}MR^{2}\right) \times \omega^{2}$$

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$$E_K = \frac{1}{4}MR^2\omega^2$$

• The combined flywheels have a total moment of inertia of

$$I_{new} = I_{1} + I_{2}$$

$$I_{new} = \frac{1}{2}MR^{2} + \frac{1}{2}\left(\frac{1}{2}M\right)\left(\frac{1}{2}R\right)^{2}$$

$$I_{new} = \frac{9}{16}MR^{2}$$

• The kinetic energy of the combined flywheels is

$$E_{K new} = \frac{1}{2} I_{new} \omega_{new}^2 = \frac{1}{2} \times \left(\frac{9}{16} MR^2\right) \times \left(\frac{2}{3}\omega\right)^2$$
$$E_{K new} = \frac{1}{2} \times \left(\frac{1}{4} MR^2\omega^2\right) = \frac{1}{2} E_K$$